

Problem Set 4

Problems to Theoretical Astrophysics, SS 2014

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1. Adiabatic exponents

a) Show that

$$\frac{\Gamma_3 - 1}{\Gamma_1} = \frac{\Gamma_2 - 1}{\Gamma_2} = \nabla_{\text{ad}}.$$

b) Determine C_V , χ_T , χ_ρ , C_P , γ , Γ_1 , Γ_2 , and Γ_3 for an ideal gas. Assume as equation of state $E = \frac{3}{2}Nk_B T$ and $PV = Nk_B T$.

c) Show the following relations for Γ_1 , Γ_2 , and Γ_3 :

$$\begin{aligned}\Gamma_1 &= \chi_\rho \gamma \\ \frac{\Gamma_2}{\Gamma_2 - 1} &= \frac{TC_P \chi_\rho}{PV \chi_T} \\ \Gamma_3 - 1 &= \frac{PV \chi_T}{T C_V}\end{aligned}$$

Hint: You can use the following formulae:

$$\begin{aligned}\gamma &= \frac{C_p}{C_v} \\ \chi_\rho &= \left(\frac{\partial \ln P}{\partial \ln \rho} \right)_T \\ \chi_T &= \left(\frac{\partial \ln P}{\partial \ln T} \right)_\rho \\ \Gamma_1 &= \left(\frac{\partial \ln P}{\partial \ln \rho} \right)_S \\ \frac{\Gamma_2}{\Gamma_2 - 1} &= \frac{1}{\nabla_{\text{ad}}} = \left(\frac{\partial \ln P}{\partial \ln T} \right)_S \\ \Gamma_3 - 1 &= \left(\frac{\partial \ln T}{\partial \ln \rho} \right)_S \\ \left(\frac{\partial x}{\partial y} \right)_z \left(\frac{\partial y}{\partial z} \right)_x \left(\frac{\partial z}{\partial x} \right)_y &= -1\end{aligned}$$

2. Schwarzschild criterion

Show that an alternate criterion for *convective stability* is given by the Schwarzschild criterion,

$$\frac{dS}{dz} > 0.$$

3. Transition to continuum and Fermi energy

- a) Assume a single particle in a box of length L . Then the momenta are quantized according to the Schrödinger equation as $\vec{p} = \frac{\hbar\pi}{L}\vec{n}$, where \vec{n} are the quantum number for each dimension. In the limit of large lengths L , discrete sums over energy levels can be transformed into integrals over the continuous momentum,

$$\frac{1}{V} \sum_{\vec{n}, \sigma} f(\epsilon_{\vec{n}}) \rightarrow g_s \int_0^\infty dp f(\epsilon(p)) v(p),$$

where $V = L^3$ is the volume, σ is the spin, and g_s the spin degeneracy.

Hint: Include an additional factor $(\Delta n)^3$, where $\Delta n = 1$, and transform it to $(\Delta p)^3$. For large L , this renders the infinitesimal quantity d^3p . Take into account $\vec{n} \geq 0$.

- b) For an ideal Fermi gas at $T = 0$, all states are occupied up to the Fermi energy ϵ_F because of the Pauli principle. The number of states equals the number of particles and is given by $N = \sum_{\vec{n}, \sigma} n_F(\epsilon)$, where $n_F(\epsilon)$ is the Fermi function.

Compute the Fermi momentum p_F and the Fermi energy ϵ_F by assuming a step function for the Fermi function.

Hint: The final result is

$$\epsilon_F = \frac{\hbar^2}{2m} \left(\frac{6\pi^2 n}{g_s} \right)^{\frac{2}{3}},$$

where $n = N/V$ denotes the number density.

- c) Compute the Fermi energy and the Fermi temperature $T_F = \epsilon_F/k_B$ for nuclear matter, a metal, and a white dwarf.

Hint: For nuclear matter, consider a nucleus obeying the relation $R = 1.25 \times 10^{-15} \text{ m } A^{1/3}$, where A is the total number of nucleons. For the other cases, the density is given by the baryonic density with $n_{\text{el}} = \rho \frac{Y_e}{m_b}$, where Y_e denotes the fraction of electrons to baryons and $m_b \approx m_u$ the mean baryonic mass. You can use the following values:

Iron	$\rho = 7.87 \frac{\text{g}}{\text{cm}^3}$	$Y_e = 0.466$
White dwarf	$\rho = 3 \times 10^6 \frac{\text{g}}{\text{cm}^3}$	$Y_e = 0.5$

4. Degeneracy in stars

A criterion for degeneracy in the interior of stars can be obtained by comparing the Fermi energy ϵ_F to the mean thermal energy $\frac{3}{2}k_B T$.

- a) Obtain a criterion for $T/\rho^{2/3}$ allowing a statement concerning the degree of degeneracy.
- b) Compare the Sun and the white dwarf Sirius B with regard to the degeneracy of electrons in the interior.

Hint: You can use the following values:

$$\begin{array}{ll} \text{Sun} & T_c = 1.58 \times 10^7 \text{K} \quad \rho_c = 162 \frac{\text{g}}{\text{cm}^3} \\ \text{Sirius B} & T_c = 7.6 \times 10^7 \text{K} \quad \rho_c = 3 \times 10^6 \frac{\text{g}}{\text{cm}^3} \end{array}$$